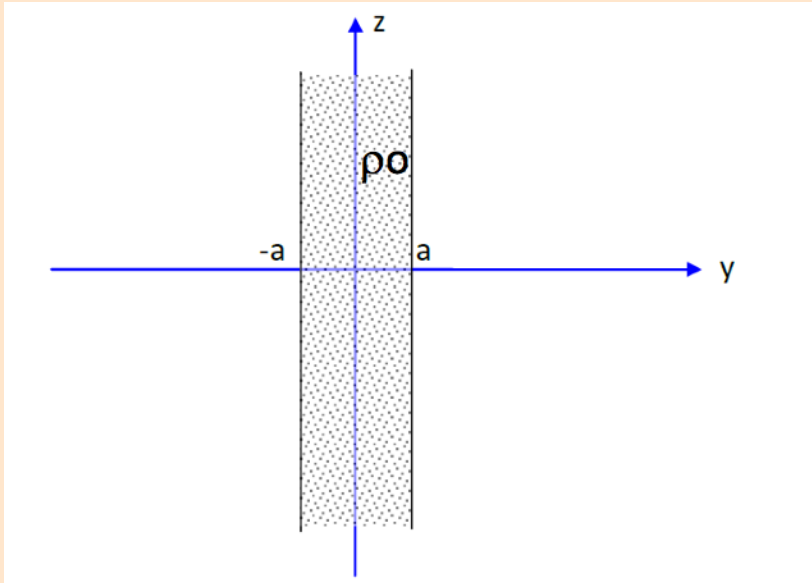


Electric Field and Potential for an infinite slab with uniform charge density ρ

PROBLEM: Consider an infinite slab parallel to the x-z plane (normal to the y axis) of uniform charge density ρ and thickness $2a$, as shown:



Find the Electric Field and Potential everywhere.

(a) First, present your arguments for treating this problem as one dimensional.

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The infinite slab can be thought of a set of parallel infinite sheets of uniform surface charge density σ ($= \rho dy$ where dy is the 'thickness of the charge sheet'). The field for such a charge sheet is independent of x and z and normal to the charge sheet, therefore normal to the x-z plane (for $\sigma > 0$, E points away from the plane of charge).

Superposition will lead to an E Field that will also be independent of x & z and normal to the x-z plane. Thus, a one D problem (y dependent).

We wish to derive the E field, $\vec{E}(y)$, at all values of y .

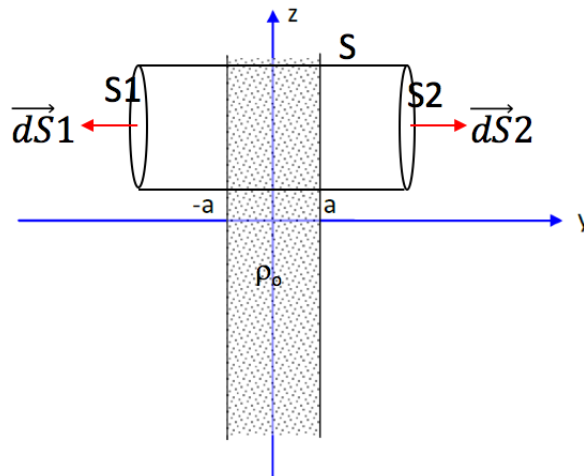
(b) Start with E outside the slab.

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Suggest using Gauss's law.

First we find the field OUTSIDE using a pillbox wider than $2a$ and both ends of area A placed OUTSIDE

of the charge slab.



\vec{E}_{outside} will be in the $\pm y$ direction and therefore normal to the ends S1 and S2 (i.e., in the $\pm \hat{y}$ direction). Let A be the area of both S1 and S2.

$$\text{Gauss's Law, } \int_S \vec{E} \cdot d\vec{S} = \int_{S1} \vec{E1} \cdot d\vec{S1} + \int_{S2} \vec{E2} \cdot d\vec{S2} = \frac{1}{\epsilon_0} (\text{Charge Contained inside } S)$$

The charge contained in the pillbox = $\rho A (2a)$.

All the E's are \parallel to the dS's; For a given y, the E's are constant.

$$E_{\text{outside}} A + E_{\text{outside}} A = \frac{1}{\epsilon_0} \rho A (2a) \text{ or}$$

$$\text{On the Right: } \vec{E}_{\text{outpos}} = \frac{1}{\epsilon_0} \rho a \hat{y}$$

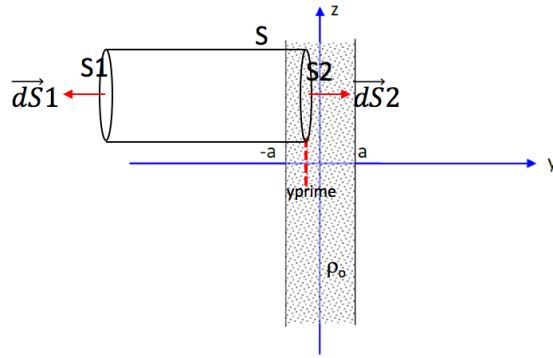
and

$$\text{On the Left: } \vec{E}_{\text{outneg}} = - \frac{1}{\epsilon_0} \rho a \hat{y}$$

(c) Now find E inside the slab.

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Now inside the slab ($-a \leq y \leq a$). We redraw S to have one end of the pillbox inside the charged region located at $y = y_{\text{prime}}$. I label the end of the pillbox inside S2 (on the right). $\vec{E1}$ at S1 = $\vec{E}_{\text{outneg}} = - \frac{1}{\epsilon_0} \rho a \hat{y}$ on the left.



Gauss's Law, $\oint_S \vec{E} \cdot d\vec{S} = \oint_{S1} \vec{E1} \cdot d\vec{S1} + \oint_{S2} \vec{E2} \cdot d\vec{S2} = \frac{1}{\epsilon_0}$ (Charge Contained inside S)

$E1 A + E_{\text{inside}}[y] A = \frac{1}{\epsilon_0}$ (Charge Contained inside S)

$\frac{1}{\epsilon_0} \rho a A + E_{\text{inside}}[y] A = \frac{1}{\epsilon_0} (\rho (a + y) A)$ Note: SIGN of y critical here in determining Q_{enclosed}

$$E_{\text{inside}}[y] = \frac{1}{\epsilon_0} (\rho (a + y) A) - \frac{1}{\epsilon_0} \rho a A$$

$$\vec{E}_{\text{inside}}[y] = \frac{\rho y}{\epsilon_0} \hat{y} \quad \text{Note: the sign of } y \text{ is critical in determining the direction of } \vec{E}_{\text{inside}}[y]$$

(d) Plot the magnitude of $E[y]$ for both inside and outside the slab.

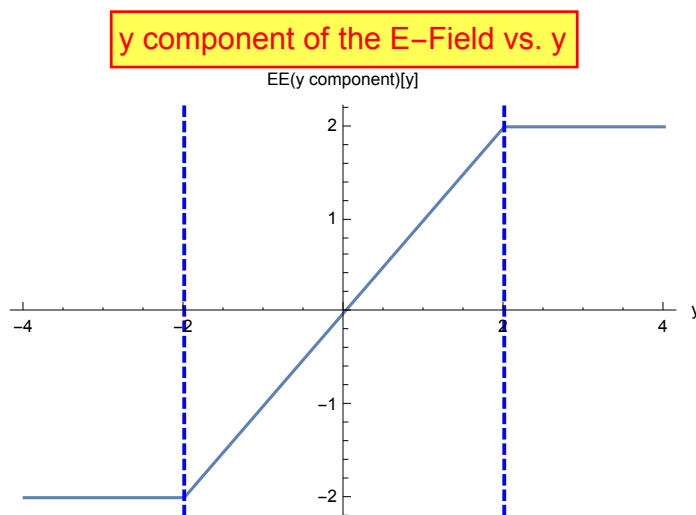
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(* Input Cell - Enter your code here -- start with a ClearAll["`*"] *)
ClearAll["`*"]

Einside[y_] =  $\frac{\rho y}{\epsilon_0}$ ; (* NOTE x takes care of the SIGN *)

Eoutneg[y_] = -  $\frac{\rho a}{\epsilon_0}$ ;

Eoutpos[y_] =  $\frac{\rho a}{\epsilon_0}$ ;

rho = 1; epsilonzero = 1; a = 2;
EE[y_] = If[y < -a, Eoutneg[y], If[-a ≤ y ≤ a, Einside[y], If[y > a, Eoutpos[y]]];
EEplot = Plot[EE[y], {y, -2 a, 2 a}, AxesLabel → {"y", "EE(y component) [y]"},
  PlotLabel → Style[Framed["y component of the E-Field vs. y"],
    16, Red, Background → Lighter[Yellow]]];
lineleft = Graphics[{Thick, Dashed, Blue, Line[{{-a, -2 a - .5}, {-a, 2 a + .5}}]}];
lineright = Graphics[{Thick, Dashed, Blue, Line[{{a, -2 a - .5}, {a, 2 a + .5}}]}];
Show[EEplot, lineleft, lineright]
```



(e) Comment on this plot

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First, the blue vertical lines are the edges of the slab of charge. Note to the left of the origin, $E[y] < 0$, meaning that the vector $\vec{E}[y]$ points to the left. To the right of the origin, $E[y] > 0$, meaning that the vector $\vec{E}[y]$ points to the right. We assumed for this plot that $\rho = +1$ (positive charge). So outside the slab \vec{E} should point away from the slab which it does. INSIDE: think of the slab as a bunch of identical, parallel sheets of positive charge. If you are at a point to the left of the origin (say $y = -1$), you would have more sheets of + charge on the right than on the left. Therefore the superposition of the fields from these sheets would lead to a field pointing left. The linear behavior of $|\vec{E}|$ is reasonable in terms of the Number of sheets on the left vs. on the right.

Not that $E[y]$ is continuous at the edges of the slab. There is a boundary condition on \vec{E} at an interface (let's use "left" and "right" for position, assuming an interface, such as ours, parallel to the x-z plane).

$$E_{\text{normal-left}} - E_{\text{normal-right}} = \frac{\sigma}{\epsilon_0}$$

For a continuous slab of charge, the surface is considered free of charge (because ρ changes abruptly from 0 to ρ with y) so we take $\sigma = 0$. Thus, $E_{\text{normal-left}}^{\text{interface}} = E_{\text{normal-right}}^{\text{interface}}$, which is what is observed in the graph at both interfaces (at $y = \pm a$).

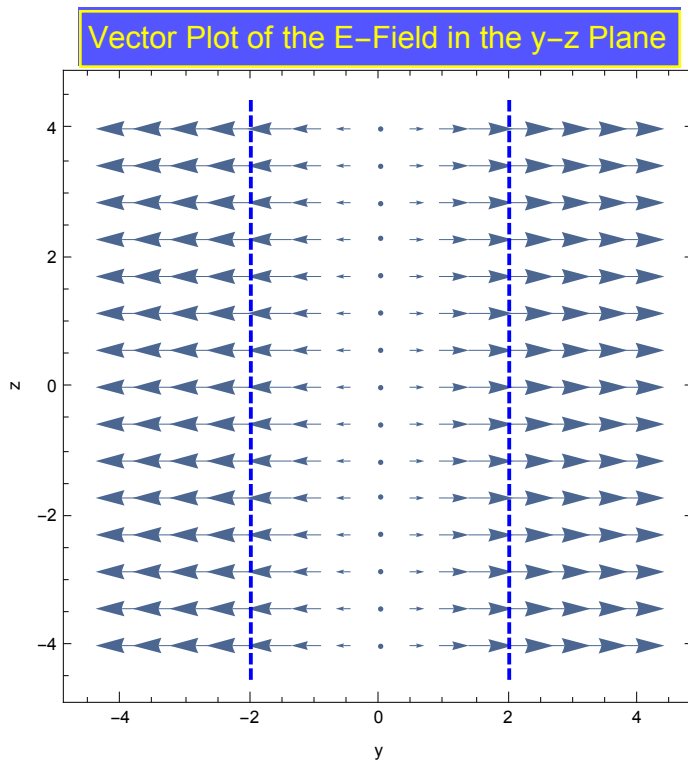
(f) Now generate a 2D Vector Plot of the E-Field both inside and outside of the slab of charge (assume a positive ρ).

(* Input Cell - Enter your code here *)

vecplot =

```
VectorPlot[{EE[y], 0}, {y, -2 a, 2 a}, {z, -2 a, 2 a}, FrameLabel -> {"y", "z"},
  PlotLabel -> Style[Framed["Vector Plot of the E-Field in the y-z Plane"],
    16, Yellow, Background -> Lighter[Blue]]];
```

Show[vecplot, lineleft, lineright]



(g) Comment on this vector plot

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Again, the blue vertical lines are the edges of the slab of charge. Again, we assumed for this plot that $\rho = +1$ (positive charge). We see that the direction of the E vectors are consistent with the amount of +

charge on the left side of a given point (y value) chosen vs. the amount of + charge on the right side of that point. $E = 0$ at $y = 0$ because the two amounts of charge balance. [Of course, we cannot forget that the slab extends to ∞ in both the x and z directions].

So E grows in magnitude (linearly with y) from zero at $y = 0$ to a maximum at $y = \pm a$. Outside, both left and right, the fields are of constant magnitude (uniform). It probably need not be mentioned, but such uniform fields require that the charge be of infinite extent in {x, z}.

Although not so obvious in the vector plot, $E_{\text{normal-left}}^{\text{interface}} = E_{\text{normal-right}}^{\text{interface}}$ at both $y = \pm a$.

(h) Derive and plot the 1D Electric Potential vs. y for the above E Field.

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Starting again with the E field expressions, we assume that the potential at $y = 0$ is zero; $\text{Potin}[0] = 0$.

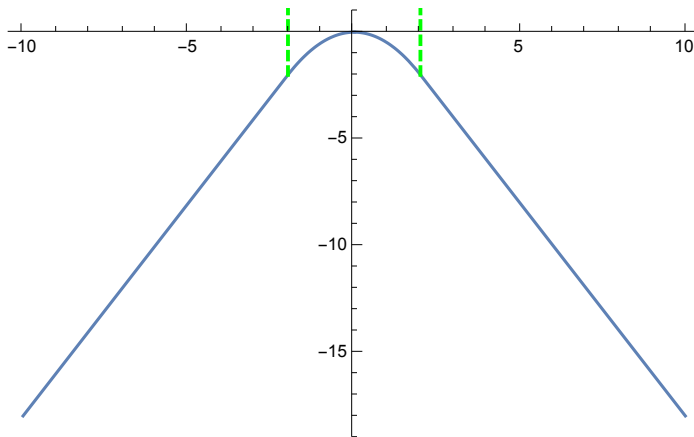
WHY? For $|E_{\text{out}}|$ constant vs. y, we would conclude that E at infinity = this constant. I.e., it would take an infinite amount of work to move a unit test charge from infinity to the slab. Thus, we could not set $\text{Potin}[\infty] = 0$. We therefore adopt $y = 0$ as a reference point where we set $\text{Potin}[0] = 0$.

NOTE: y takes care of SIGNS; we use yprime as the variable of integration:

```
(* Input Cell - Enter your code here *)
ClearAll["`*"]
Ein[y_] =  $\frac{\rho y}{\epsilon_0}$ ;
Eoutneg[y_] = -  $\frac{\rho a}{\epsilon_0}$ ;
Eoutpos[y_] =  $\frac{\rho a}{\epsilon_0}$ ;
Potin[y_] = -  $\int_0^y \text{Ein}[y_{\text{prime}}] \, dy_{\text{prime}}$ 
Potoutpositive[y_] = Potin[a] -  $\int_a^y \text{Eoutpos}[y_{\text{prime}}] \, dy_{\text{prime}}$ 
Potoutnegative[y_] = Potin[-a] -  $\int_{-a}^y \text{Eoutneg}[y_{\text{prime}}] \, dy_{\text{prime}}$ 
Pottot[y_] =
  If[y < -a, Potoutnegative[y], If[-a ≤ y ≤ a, Potin[y], Potoutpositive[y]]];
rho = 1; epsilonzero = 1; a = 2;
PotPlot = Plot[Pottot[y], {y, -5 a, 5 a}];
lineleft = Graphics[{Thick, Dashed, Green, Line[{{-a, -a}, {-a, a}}]}];
lineright = Graphics[{Thick, Dashed, Green, Line[{{a, -a}, {a, a}}]}];
Show[PotPlot, lineleft, lineright]
```

$$-\frac{\rho y^2}{2 \epsilon_0}$$

$$-\frac{a^2 \rho}{2 \epsilon_0} - \frac{a \rho (-a + y)}{\epsilon_0}$$

$$-\frac{a^2 \rho}{2 \epsilon_0} + \frac{a \rho (a + y)}{\epsilon_0}$$


(i) Comment on this result (the plot).

<This is a text cell - enter your answer here.>

The green dashed lines show the values of y ($\pm a$) at the edges of the slab.

The straight lines correspond to the constant E fields outside of the slab with slopes indicating the magnitudes of E

$$(\vec{E}[y] = - \text{Grad}[\text{Potential}] = - \frac{dV}{dy} \hat{y}).$$

Inside the slab, E changes linearly with y so that the potential $\sim -y^2$; this is consistent with the \vec{E} pointing away from the center of the slab. Check yourself out on the SIGNS for the above and the directions of \vec{E} in the various regions.